

Physics 114 - January 31, 2006

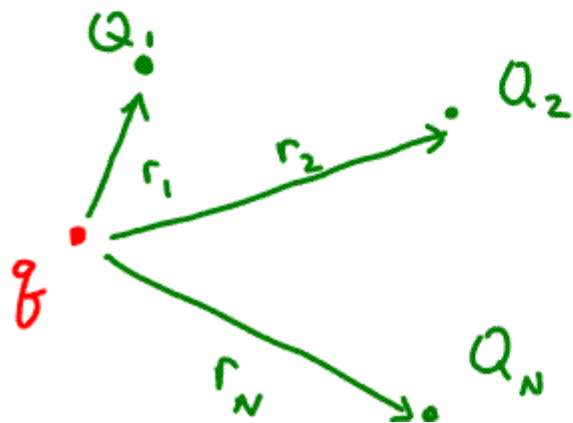
Last Time :

Charge is Conserved

Conductors

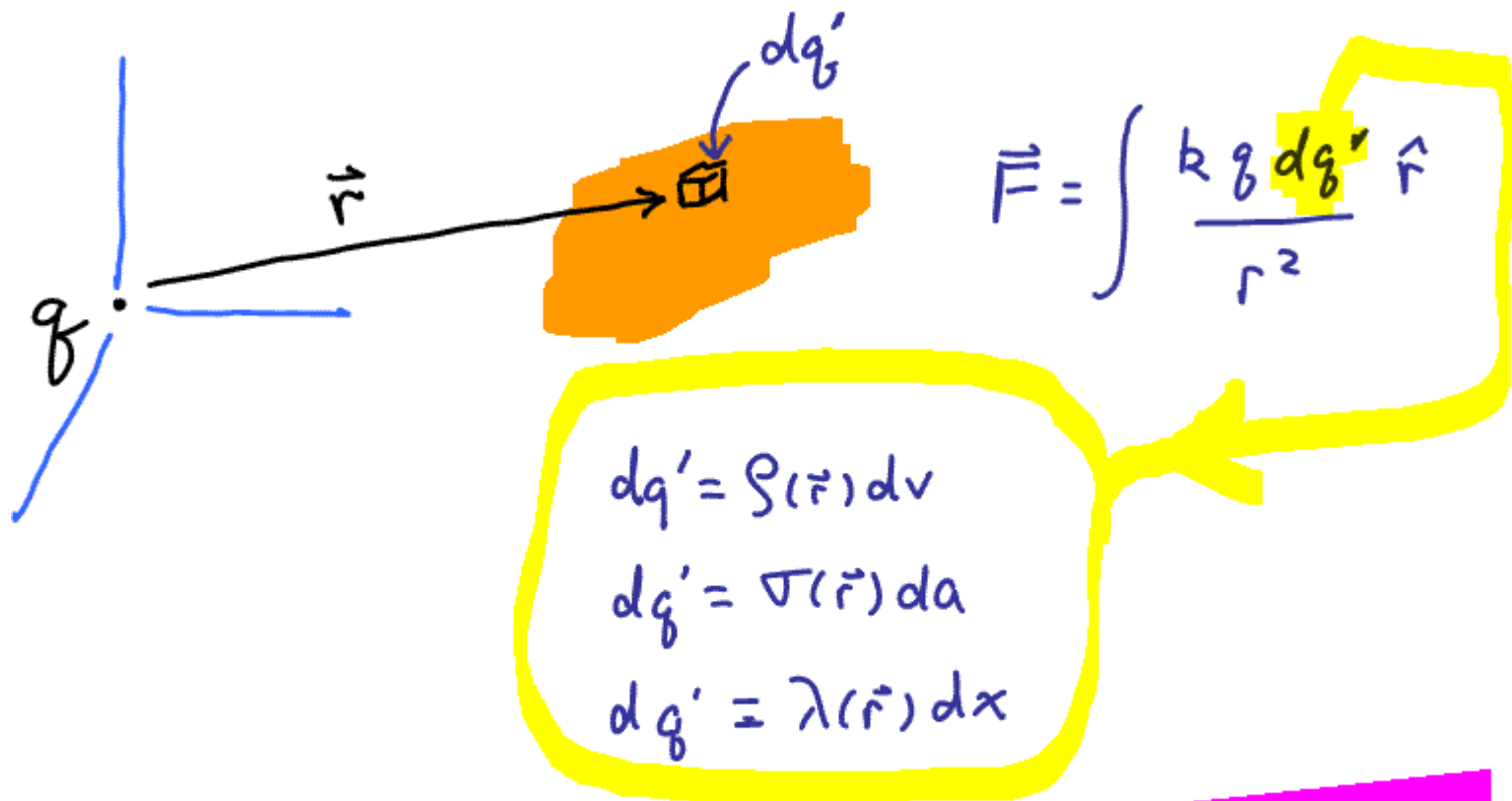
insulators

Coulomb's Law for system of discrete charges



$$\vec{F} = \sum_{i=1}^N k q Q_i \frac{\hat{r}_i}{r_i^2}$$

Coulomb's Law for a continuous charge dist.

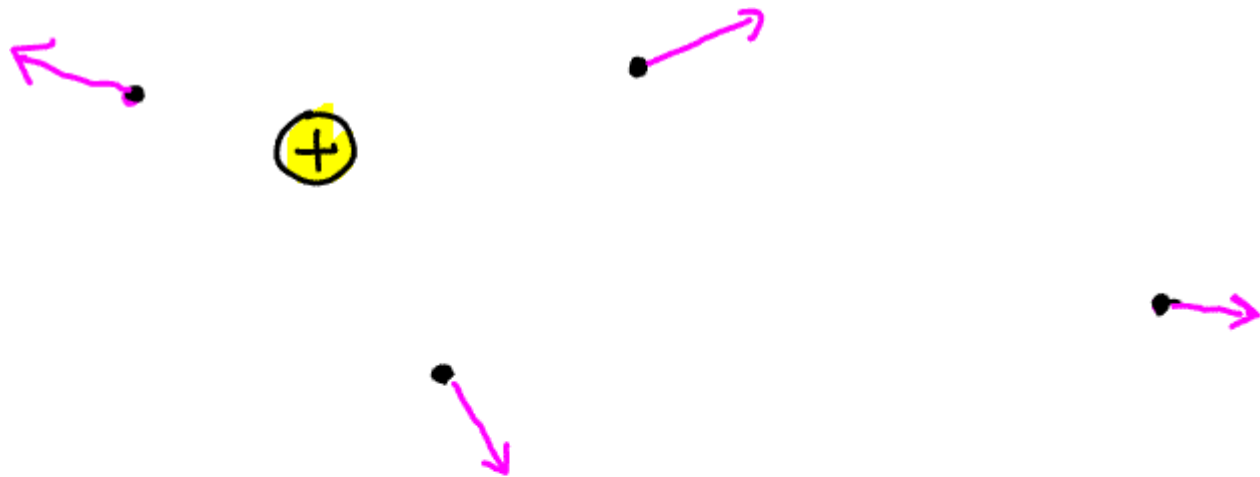


doable ... Sometimes nasty

usually work w/

Electric field instead

$$\vec{E}(\vec{r}) = \frac{\vec{F}(\vec{r})}{q}$$



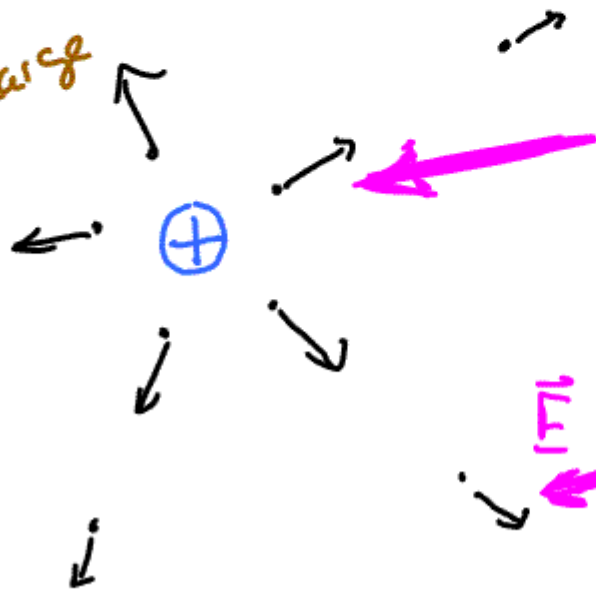
$$\vec{F} = \frac{k q_1 q_2}{r^2} \hat{r}$$

$$\vec{E} = \frac{\vec{F}}{q_1} = \frac{k q_2}{r^2} \hat{r}$$

$$\vec{E} = \vec{F}/q$$


\vec{E} vectors point radially outward
(as does force)

\vec{E} vectors
About a
⊕ Point charge

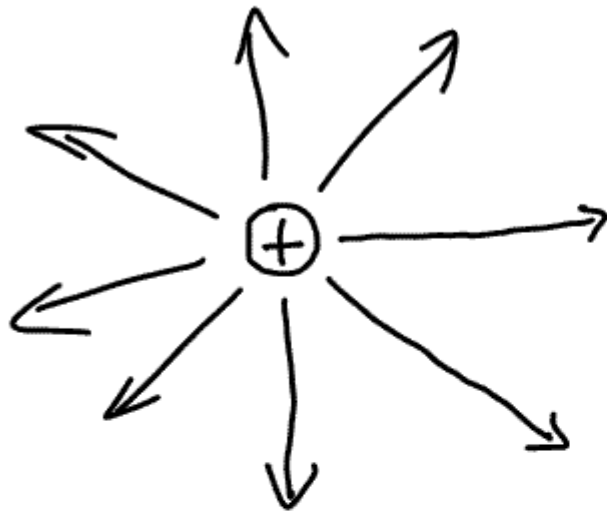


Closer ... \vec{E} vector
long ... Strong
field

Farther Away ...

\vec{E} vector shorter here
because field is
weaker ...
less force/charge

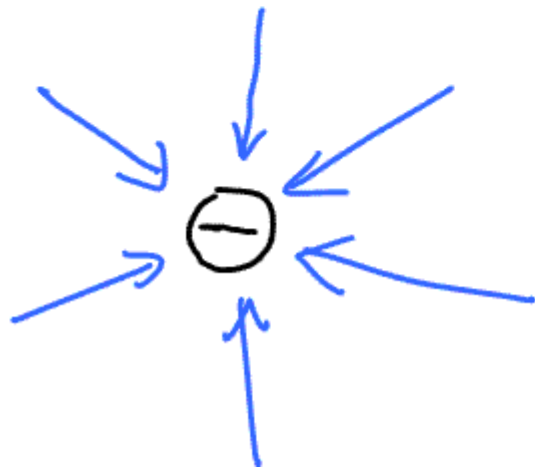
Lines of force

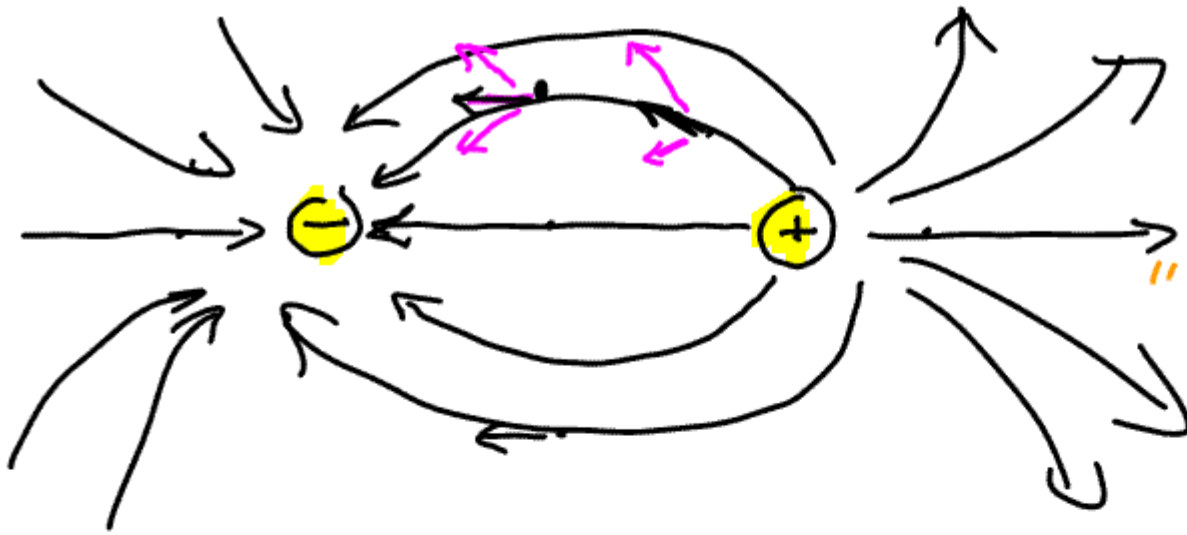


- Go in direction of + Test chg.

- density of lines
→ Strength of force/field

- Never cross

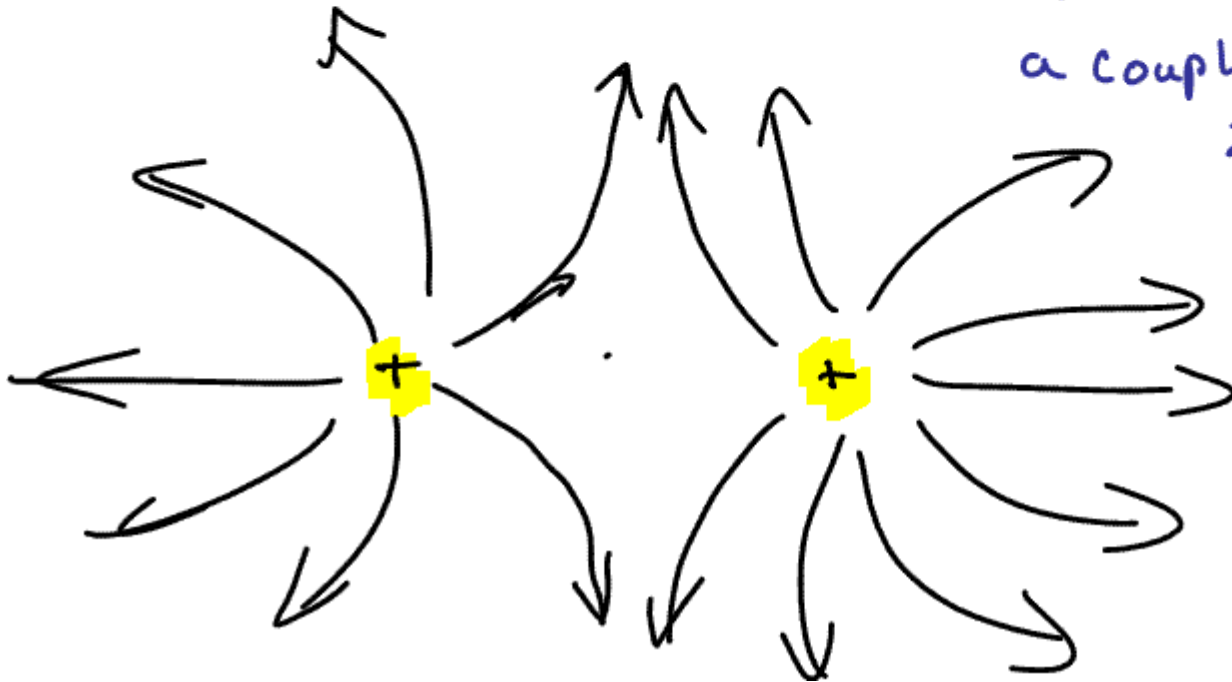


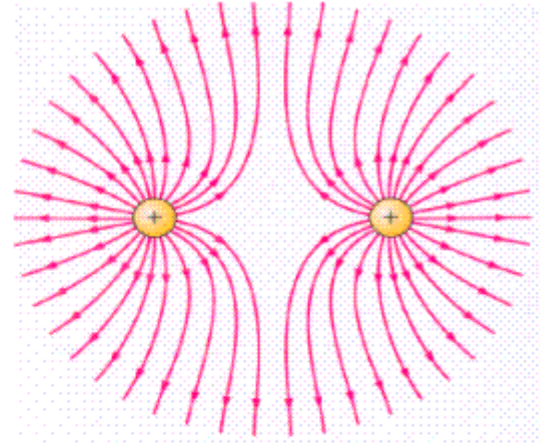
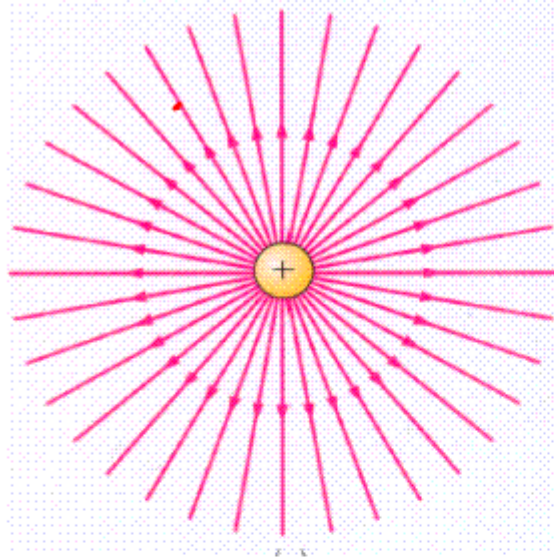


"Electric Dipole Field"

Lines of force around a couple of 2-charge

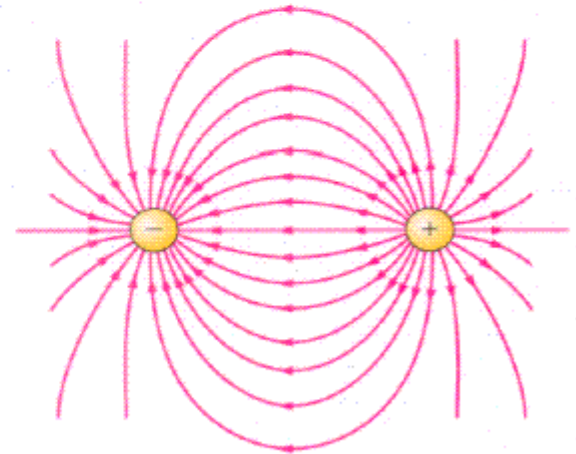
Distributions





Lines of force

Nicer figures
from textbook
by Tipler



Lines of force are a useful
way to visualize electric fields

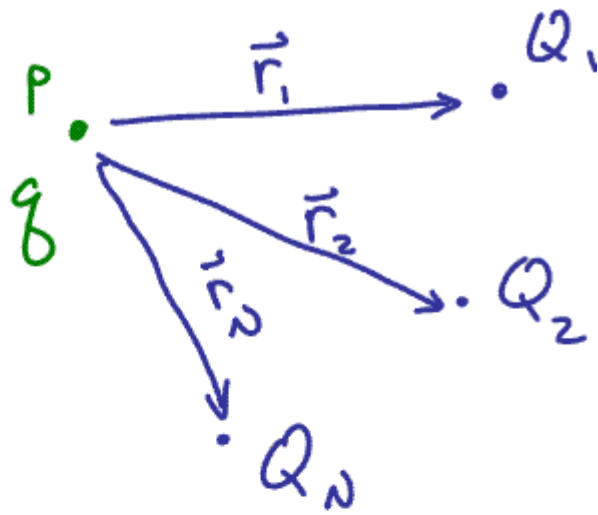
\vec{E} for a point charge



$$\vec{F}_{\text{at } P \text{ due to } Q} = \frac{k Q q_0}{r^2} \hat{r}$$

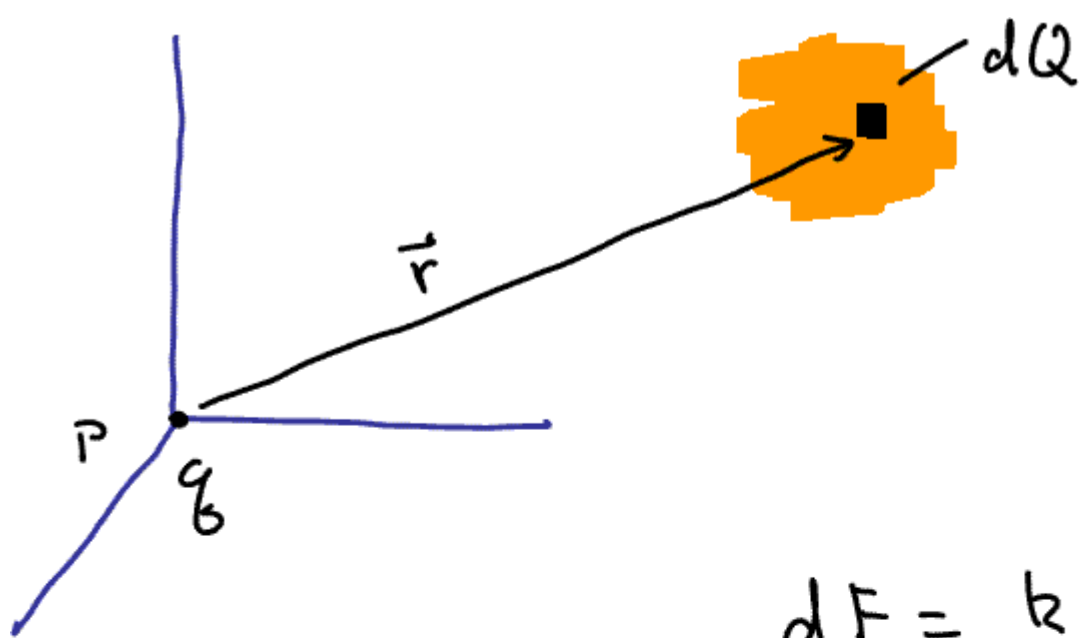
$$\vec{E} = \frac{\vec{F}}{q_0}$$

$$\vec{E}_{\text{at } P \text{ due to } Q} = \frac{k Q}{r^2} \hat{r}$$



$$F_{at P} = \sum_{i=1}^N \frac{k Q_i g}{r_i^2} \hat{r}_i$$

$$F_{at P} = \sum_{i=1}^N \frac{k Q_i}{r_i^2} \hat{r}_i$$



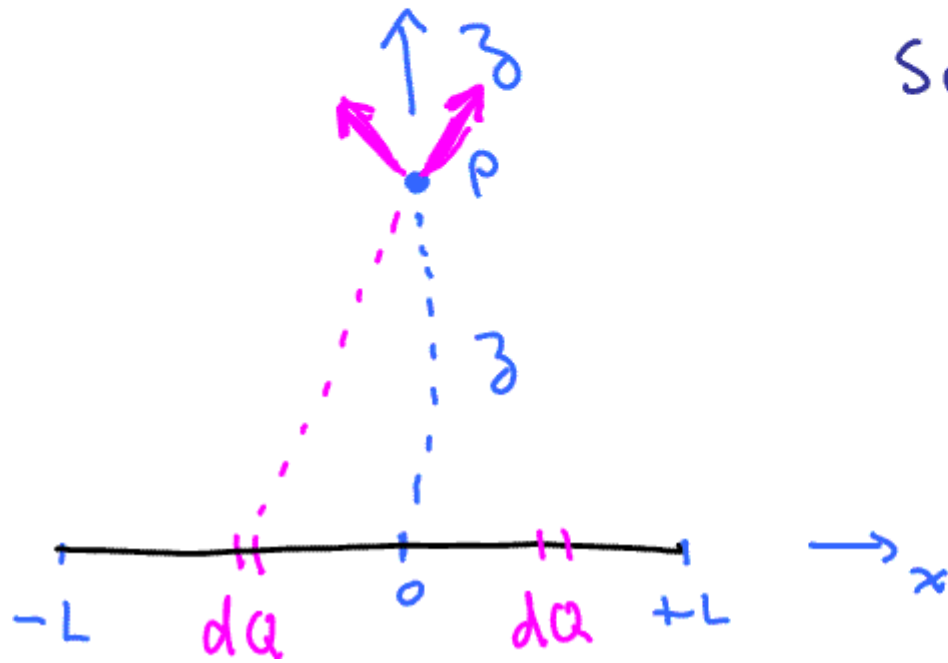
$$dF = \frac{k q_b dQ}{r^2} \hat{r}$$

$$d\vec{E} = \frac{k dQ}{r^2} \hat{r}$$

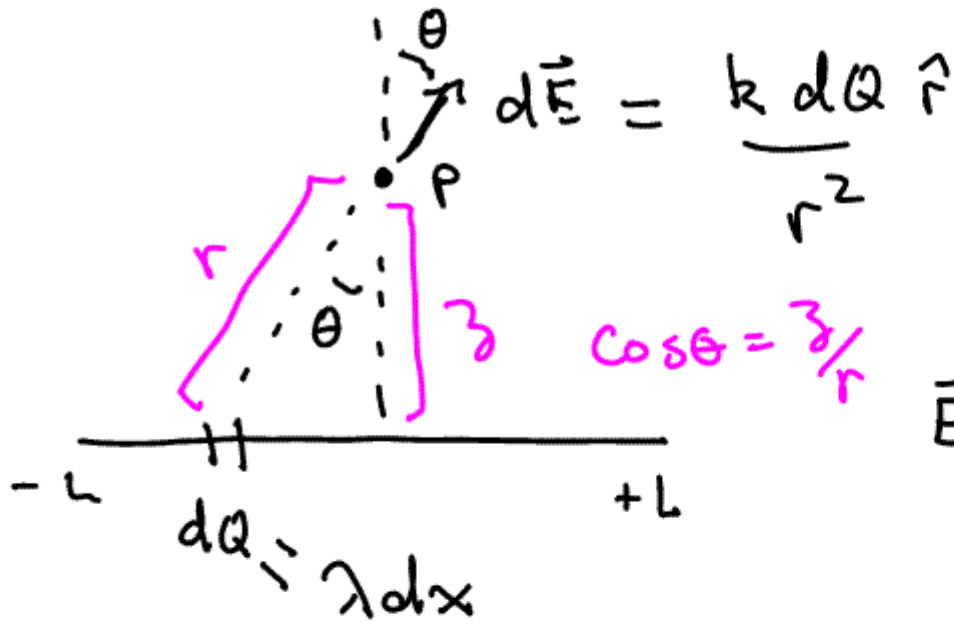
$$\vec{E}_{\text{at } P} = \int \frac{k dQ}{r^2} \hat{r}$$

chg
dist

Find \vec{E} at a distance z above the midpoint of a line segment of length $2L$ that carries a uniform line charge of $+\lambda$



Symmetry
 $\rightarrow \vec{E}$ STRAIGHT UP



$$d\vec{E} = \frac{k dQ}{r^2} \hat{r}$$

$$\cos\theta = \frac{z}{r}$$

$$\vec{E} = \int_{-L}^{+L} \frac{k dQ}{r^2} \hat{r}$$

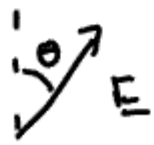
$$\vec{E} = \int_{-L}^{+L} \frac{k \lambda dx}{r^2} \hat{r}$$

Symmetry

only worry abt

straight up component

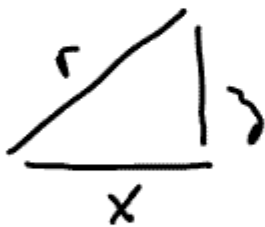
$E \cos\theta$



in \hat{z} direction

$$E_{\text{net}} = \int_{-L}^L \frac{k\lambda}{r^2} dx \quad \cancel{r} \quad \cos\theta \quad \hat{u}_p \quad \text{or} \quad \hat{z}$$

$$E_{\text{net}} = \int_{-L}^L \frac{k\lambda}{r^2} dx \frac{z}{r} \hat{z}$$



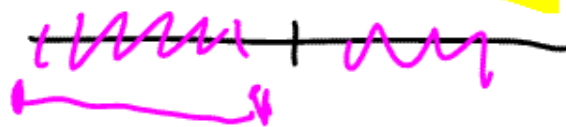
$$r = (x^2 + z^2)^{1/2}$$

$$E_{\text{net}} = \int_{-L}^L \frac{k\lambda dx}{(x^2 + z^2)^{3/2}} z \hat{z}$$

By symmetry of problem
change limits from
 $-L \rightarrow L$ to $0 \rightarrow L$
and mult
by 2

$$E_{\text{net}} = k\lambda z \int_{-L}^L \frac{dx}{(x^2 + z^2)^{3/2}} = 2k\lambda z \int_0^L \frac{dx}{(x^2 + z^2)^{3/2}} \hat{z}$$

check units
check limiting
cases



$$\vec{E} = \frac{2k\lambda L}{z(L^2 + z^2)^{1/2}} \hat{z}$$

$$\vec{E} = \frac{F}{q} = \frac{\text{Newton}}{\text{Coul}}$$

$$\frac{kq^2}{r^2} \sim \text{Newton}$$

$z \rightarrow \text{large}$

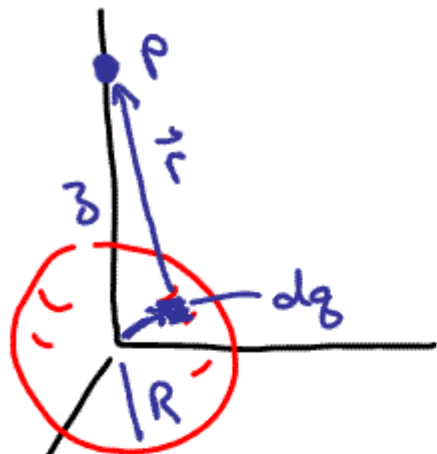
$$\frac{2k\lambda}{z(L^2 + z^2)^{1/2}} \rightarrow \frac{2k\lambda L}{z^2}$$

charge

Field of a point charge with
total charge λL viewed
from a distance z

✓
Limiting
case
Makes sense

Sphere of
uniform
charge



$$\rho = \frac{Q_{\text{TOT}}}{\frac{4}{3}\pi R^3}$$

Wally

$$\vec{E} = \int_{V_{01}} \frac{k(dq)}{r^2} \hat{r} = \rho dV$$

Crash

$$= \int_{V_{01}} \frac{k \rho dV}{r^2} \hat{r} = \int_0^R \int_0^\pi \int_0^{2\pi} \frac{k \rho}{r^2} r^2 \sin \theta d\theta d\phi dr$$

This can be quite difficult. There are easier ways. I drag you through the mud here so you can appreciate the usefulness of Gauss' Law (Next Time)