

# Interactions of charged particles w/ Matter

⇒ ionization (lower energies)

Show  $dE/dx$  curve

more of  
higher  
Mass!

⇒ Radiation - Bremsstrahlung

$$\left(\frac{dE}{dx}\right)_{\text{rad}} \propto E$$

whereas

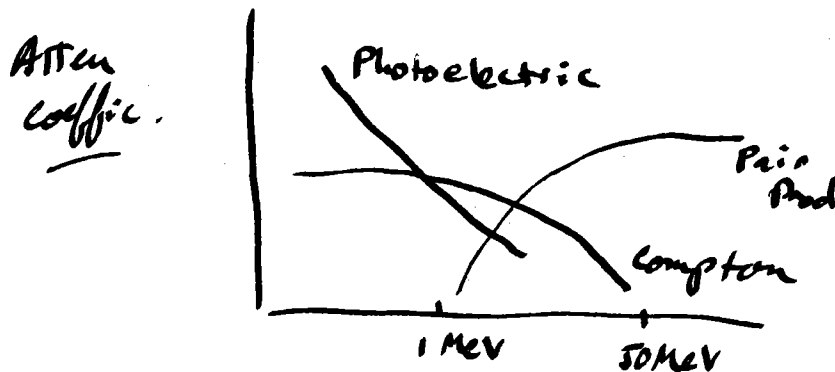
$$\left(\frac{dE}{dx}\right)_{\text{ion}} \sim \text{constant}$$

Bio damage from ionization

& worse than  $\beta$  (outside body)

Also penetrates less far

## Interactions of $\gamma$ w/ matter



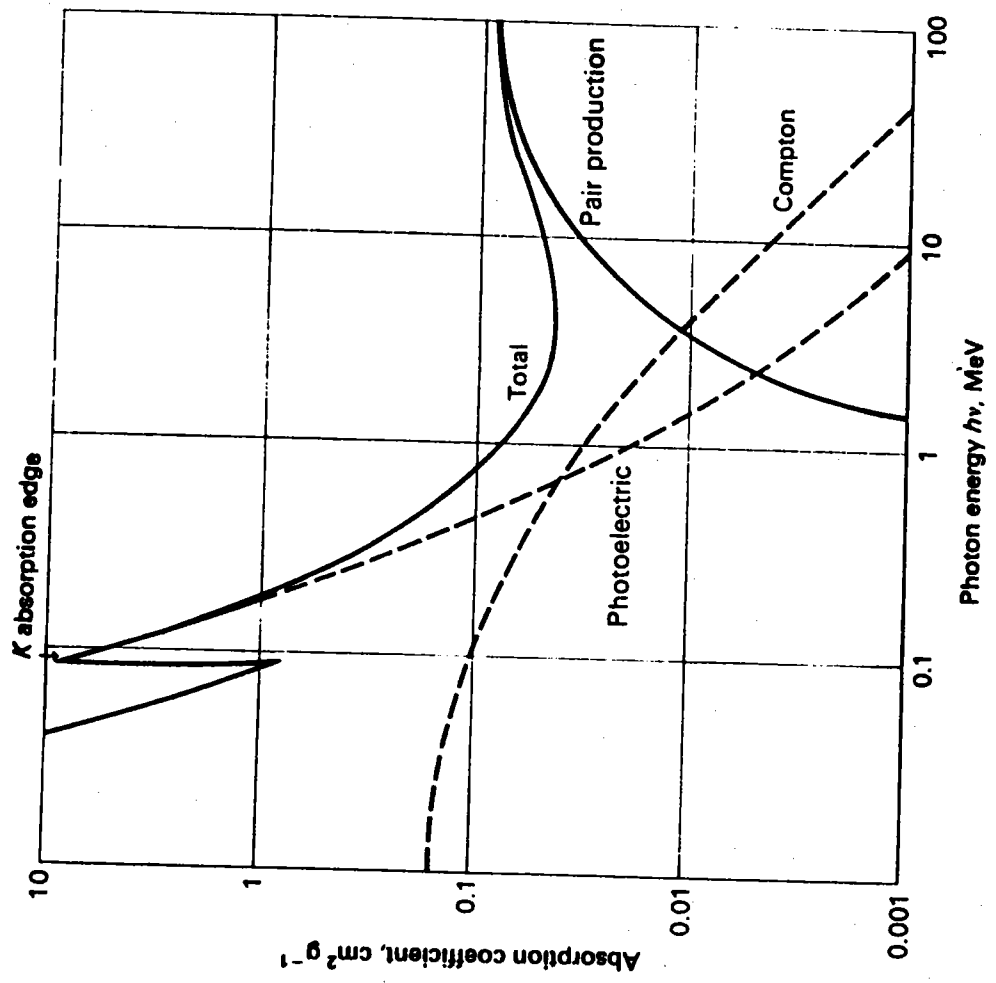
$$I = I_0 e^{-\mu x}$$

$\frac{I}{I_0} = \text{Conversion length}$

$$\frac{I}{I_0} = e^{-\mu x}$$

$$\mu = I_0 / e$$

$\mu$



the thickness of absorber because of the range. The energy loss

is given by Eq. 2.12, the range is given

$$\frac{dE}{M} = \frac{dE}{M} \quad (2.13)$$

of  $E/M$ . To illustrate the usefulness of Eq. 2.13, consider a proton, for example, has a range of  $R$  in a material of density  $\rho$  and atomic number  $Z$ . Then the range of another particle, for example, with energy  $E_a$  can be related to

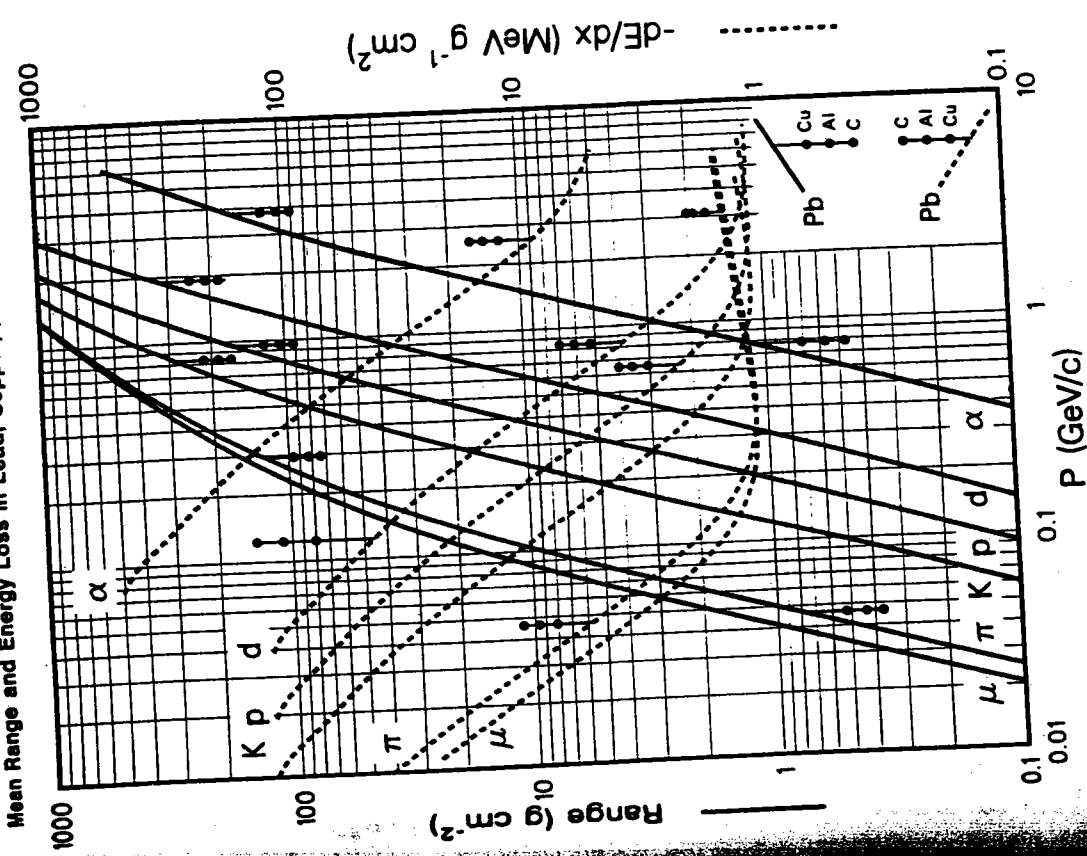
$$\left(\frac{E_a}{M_a}\right) R_a = \left(\frac{E_p}{M_p}\right) R_p \quad (2.14)$$

an often be expressed empirically in the form of Eq. 2.14. For low energy protons in air can be taken as  $R_p = 9.3 \text{ MeV} [3]$ . The range and  $dE/dx$  curves for a number of incident particles in a material of mass traversed ( $\text{g}/\text{cm}^2$ ) instead of  $dE/dx$  curves show the  $1/v^2$  drop for small momentum ionization for higher momentum. Particles in liquid hydrogen in Fig. 2.4.

the energy loss of the energy loss given in the present nature of the particles in an ad hoc must take into account (1) the fact that

Figure 2.3 Mean range and energy loss of charged particles in solids. Calculations use the Bethe-Bloch equation with density effect corrections. Refer to the cited reference for a discussion of assumptions and qualifications. (Particle Data Group, Rev. Mod. Phys. 56: S1, 1984.)

**PARTICLE DETECTORS, ABSORBERS, AND RANGES**  
**Mean Range and Energy Loss in Lead, Copper, Aluminum, and Carbon**



So one can have Nuclear Spectroscopy

i.e., Transitions of Nucleons between allowed energy states within nuclei

- EM radiation emitted/Absorbed  $\Rightarrow$   $\gamma$  rays
- Energies between states much larger than for Atoms

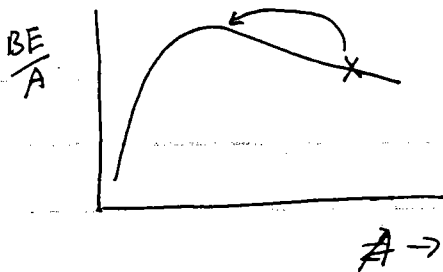
Emitted EM Radiation  $\gamma$  rays ... Very high energy EM rad.

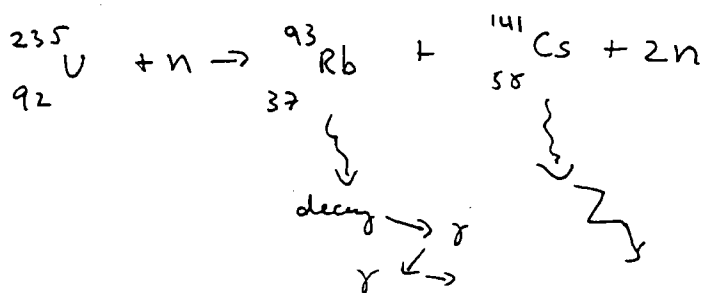
$\alpha < \beta < \gamma$  Bio damage outside body

$\alpha > \beta > \gamma$  Bio damage if substance incorporated in body  
 (This is a fn of where/if substance is concentrated in some organ also)  
 (also a fn of substance activity)

### Nuclear Fission

useful for winning wars & heating tea





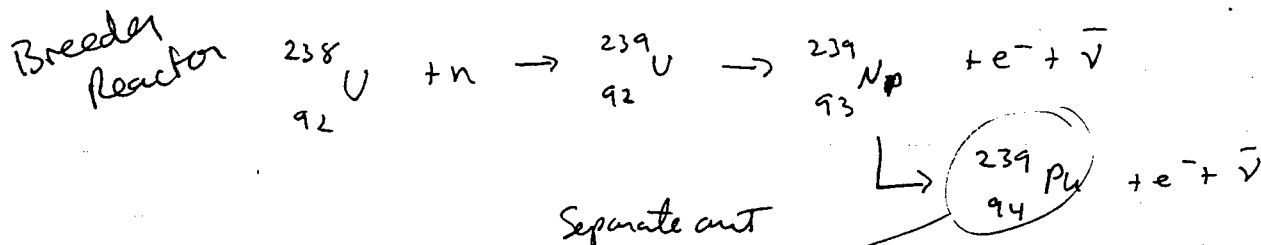
~ 200 MeV of energy released (KE +  $\gamma$  rays)  
 + 2 n to start new fissions

uncontrolled  $\rightarrow$  bomb

Add moderator to stop neutrons  $\rightarrow$  can control reaction  
 get a reactor

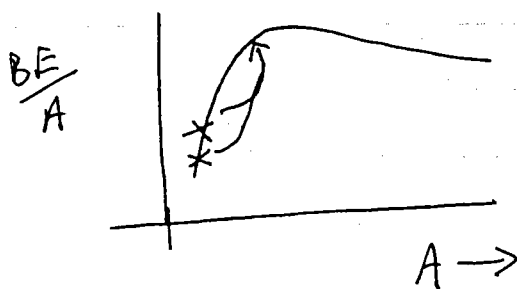
# neutrons available to start new fission rxns = 1  $\Rightarrow$  Critical  
 $\gg 1$  Bomb  
 $\ll 1$  No rxn

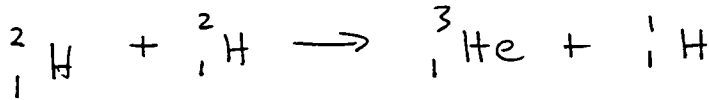
Fortunately  ${}^{235}\text{U} \sim 0.7$  percent  
 ${}^{238}\text{U} \sim 99.3$  percent } Hard to isolate  ${}^{235}\text{U}$



can be used for nuclear fuel + bombs

## Nuclear Fusion





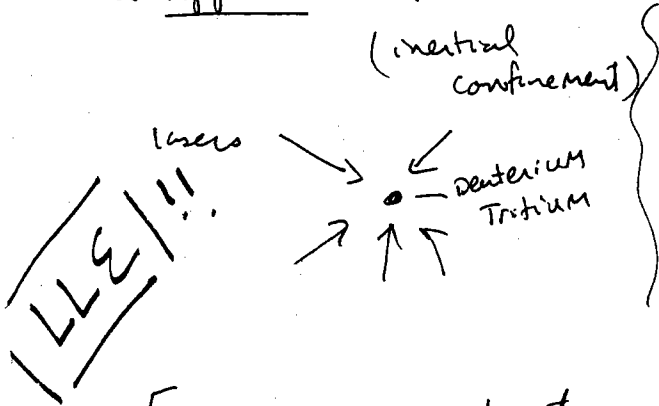
liberates  $\sim 1 \text{ MeV/nucleon}$  in energy

- get H from water  $\Rightarrow$   $\infty$  supply
- byproducts are "clean"

problem: Must overcome Coulomb repulsion

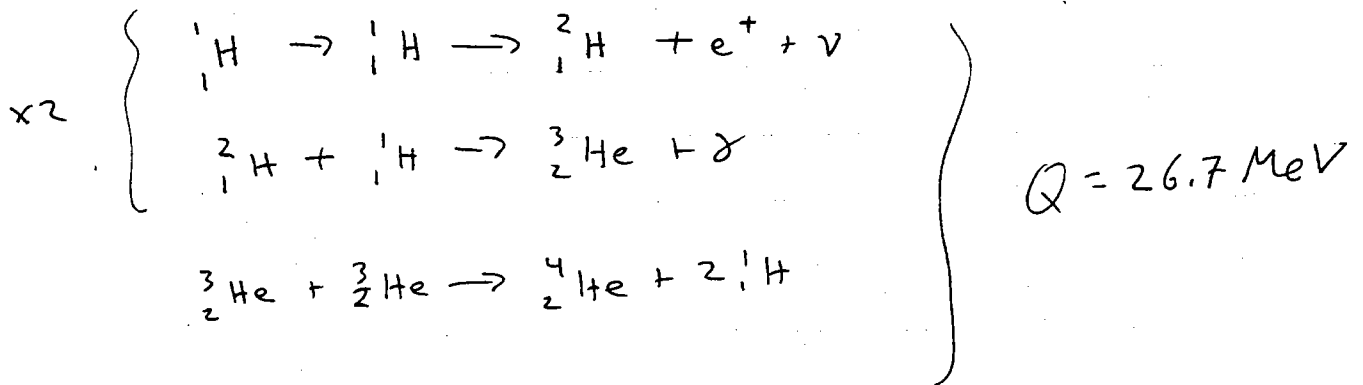
Hard to do on large scale in a controlled self sustaining manner

2 Approaches:

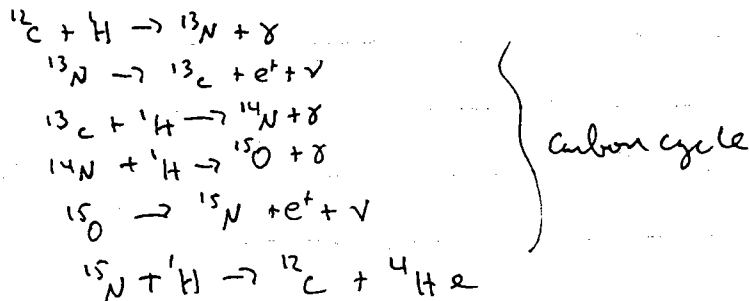


Tokamak (Magnetic confinement)  
Heat  
Plasma  
until thermal energy sufficient to ignite fuel.

Fusion is important ... Stars work by Nuclear fusion i.e. our sun



Stellar life cycle if Time Allows





radioactive

Sample of  $N_1$  nuclei of certain type  $x$ 

$$\frac{dN}{dt} \propto N$$

$$A \equiv \frac{dN}{dt} = -\lambda N$$

↑ decay constant (units  $\frac{1}{s}$ )

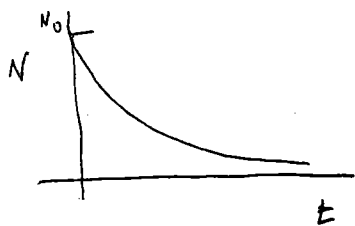
Activity

≡ change in # radioactive nuclei/s

$$\frac{dN}{N} = -\lambda dt$$

$$\ln N = -\lambda t + c$$

$$\boxed{N = N_0 e^{-\lambda t}}$$

Exponential law of  
radioactive decaycannot easily measure  $N$  so

$$\boxed{A = A_0 e^{-\lambda t}}$$

Activity measured in curies

$$1 \text{ Curie (Ci)} = 3.7 \times 10^{10} \text{ decays/s}$$

usually work in mCi +  $\mu$ Ci

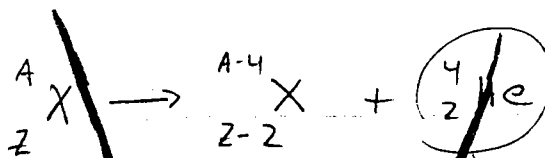
Half life  $\equiv$  time for  $A$  for a sample to be reduced by  
a factor of two

$$a = a_0 e^{-\lambda t}$$

$$\frac{1}{2} a_0 = a_0 e^{-\lambda t_{1/2}}$$

$$t_{1/2} = \frac{1}{\lambda} \ln 2 = \frac{0.693}{\lambda}$$

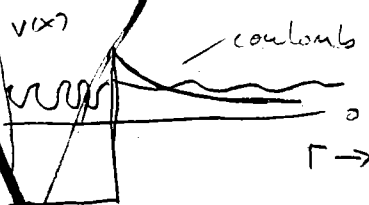
α-decay



α-particle

Protons  
Nucleons exist bouncing around inside the "bag" of the nucleus

Nuclear potential



Quantum Tunneling occurs ... α particle escapes from nucleus

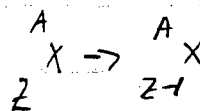
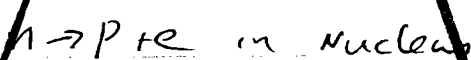
α-particle Heavy ... charge of +2

Heavily ionizing  
Easily stopped

do in several  
pages

least harmful radiation if outside body!  
Most harmful radiation if ingested!  
(given a certain activity & position in body)

β-decay



"Q-value" = KE given to products =  $(m_n - m_p - m_e) c^2$

Assume no/little nuclear recoil

expect a spike in β decay spectrum